

# Para Hamiltonian Equations With Poisson Brackets

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## Abstract

In this study, taking care of the framework of para-Kählerian manifolds it was introduced para-complex analogue of Hamiltonian equations with Poisson bracket. Also, it was shown that a canonical transformation preserves para-complex Hamiltonian equations with Poisson structure. Finally, the geometrical and mechanical conclusions on the para-mechanic systems have been suggested.

**Keywords:** *para-Kählerian manifolds; Hamiltonian systems; Poisson manifold; canonical transformation*

## Özet

Bu çalışmada, para-Kähler manifoldların çatısı dikkate alınarak, Poisson parantezli Hamilton denklemlerinin para-kompleks benzeri tanıtıldı. Aynı zamanda; bir kanonik dönüşümün Poisson yapılı para-kompleks Hamilton denklemleri koruduğu gösterildi. Sonuçta ise; para-mekanik sistemler üzerindeki geometrik ve mekanik sonuçlar tartışıldı.

**Anahtar Kelimeler:** *para- Kähler manifoldlar; Hamilton sistemler; Poisson manifold; kanonik dönüşüm*

## 1. Introduction

Differential geometry provides a good framework for studying Hamiltonian formalisms of classical mechanics. It is possible to show in [1,2,3] some numerous articles and books where differential geometric methods in mechanics are presented. In

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fact, we may say that the role of symplectic geometry in Hamiltonian theories is similar to that of tangent geometry in Lagrangian theories [3]. The dynamics of Hamiltonian formalisms is characterized by a suitable vector field defined on cotangent bundles which are phase-spaces of momentum of a given configuration manifold  $Q$  of dimension  $m$ . If  $H: T^*Q \rightarrow \mathbf{R}$  is a regular Hamiltonian function then there is a unique vector field  $Z_H$  on  $T^*Q$  such that dynamical equations

$$i_{Z_H} \Phi = \Phi(Z_H) = dH, \quad (1)$$

where  $\Phi$  is the symplectic form and  $H$  stands for Hamiltonian function. The paths of the Hamiltonian vector field  $Z_H$  are the solutions of the Hamiltonian equations. The triple  $(T^*Q, \Phi, Z_H)$  is called Hamiltonian system on the cotangent bundle  $T^*Q$  with symplectic form  $\Phi$ . Hamilton equations with Poisson structure are written as

$$dq^i/dt = \{q^i, H\}, \quad dp_i/dt = \{p_i, H\}, \quad (2)$$

where  $(q^i, p_i)$ ,  $1 \leq i \leq m$  are canonical coordinates on  $T^*Q$ .

Complex (para-complex) analogues of the Hamiltonian equations were obtained in the framework of Kählerian (para-Kählerian) manifolds and the geometric results on a complex (para-complex) mechanical systems were found [4,5]. Also, complex version of Hamiltonian equations with Poisson structure was introduced [6].

The goal of this study is to make a contribution to the modern development of Hamiltonian formalisms of classical mechanics in terms of differential-geometric methods on differentiable manifolds. From this point of view, this manuscript presents the para-complex analogues of Hamiltonian equations with Poisson bracket and discusses geometrical and mechanical conclusions on a para-mechanic systems.

The present paper is structured as follows. In section 2, it is recalled para-complex and para-Kählerian manifolds, and also para-complex analogues of Hamiltonian equations. In section 3, Poisson structure is generalized to para-Kählerian manifolds. In section 4, it is obtained para-complex version of Hamiltonian equations with Poisson bracket. In the conclusion section, geometrical and mechanical conclusions on Hamiltonian mechanics systems were suggested.

## 2. Preliminaries

In this study, all manifolds and geometric objects are differentiable and the Einstein summation convention is in use. Also, it is denoted by  $\mathbf{A}$  the set of para-complex numbers, by  $F(TM)$  the set of para-complex functions on  $TM$ , by  $\mathfrak{X}(TM)$  the set of

para-complex vector fields on  $\mathbf{TM}$  and by  $\Lambda^1(\mathbf{TM})$  the set of para-complex 1-forms on  $\mathbf{TM}$ . The definitions and geometric structures on the differential manifold  $\mathbf{M}$  given by [7] may be extended to  $\mathbf{TM}$  as follows:

### 2.1 Para-complex Manifolds

**Definition 1:** A tensor field  $\mathbf{J}$  of type (1,1) on  $\mathbf{TM}$  such that  $\mathbf{J}^2=\mathbf{I}$  is called an almost product structure on a tangent bundle  $\mathbf{TM}$  of configuration manifold  $\mathbf{M}$  of real dimension  $m$ . The pair  $(\mathbf{TM},\mathbf{J})$  is said to be an almost product manifold. An almost para-complex manifold is an almost product manifold  $(\mathbf{TM},\mathbf{J})$  such that the two eigenbundles  $\mathbf{TT}^+\mathbf{M}$  and  $\mathbf{TT}^-\mathbf{M}$  associated to the eigenvalues  $+1$  and  $-1$  of  $\mathbf{J}$ , respectively have the same rank. The dimension of an almost para-complex manifold is necessarily even. Equivalently, a splitting of the tangent bundle  $\mathbf{TTM}$  of tangent bundle  $\mathbf{TM}$ , into the Whitney sum of two sub bundles on  $\mathbf{TT}^{\pm}\mathbf{M}$  of the same fibre dimension is called an almost para-complex structure on  $\mathbf{TM}$ . An almost para-complex structure on a  $2m$ -dimensional manifold  $\mathbf{TM}$  may alternatively be defined as a  $\mathbf{G}$ - structure on  $\mathbf{TM}$  with structural group  $\mathbf{GL}(n,\mathbf{R})\times\mathbf{GL}(n,\mathbf{R})$ .

A para-complex manifold is an almost para-complex manifold  $(\mathbf{TM},\mathbf{J})$  such that the  $\mathbf{G}$ - structure defined by the tensor field  $\mathbf{J}$  is integrable. Assume that  $\mathbf{x}^i$  and  $(\mathbf{x}^i,\mathbf{y}^i)$ ,  $1\leq i\leq m$  are a real coordinate system on neighbourhoods  $\mathbf{U}_p$  and  $\mathbf{TU}_p$  of any points  $p$  and  $\mathbf{T}_p$  of  $\mathbf{M}$  and  $\mathbf{TM}$  and, also  $\{(\partial/\partial\mathbf{x}^i)_p, (\partial/\partial\mathbf{y}^i)_p\}$  and  $\{(\mathbf{d}\mathbf{x}^i)_p, (\mathbf{d}\mathbf{y}^i)_p\}$  natural bases over  $\mathbf{R}$  of the tangent space  $\mathbf{T}_p(\mathbf{TM})$  and the cotangent space  $\mathbf{T}_p^*(\mathbf{TM})$  of  $\mathbf{TM}$ , respectively. It can be seen to be

$$\mathbf{J}(\partial/\partial\mathbf{x}^i)=\partial/\partial\mathbf{y}^i, \mathbf{J}(\partial/\partial\mathbf{y}^i)=\partial/\partial\mathbf{x}^i \quad (3)$$

and

$$\mathbf{J}^*(\mathbf{d}\mathbf{x}^i)=-\mathbf{d}\mathbf{y}^i, \mathbf{J}^*(\mathbf{d}\mathbf{y}^i)=-\mathbf{d}\mathbf{x}^i. \quad (4)$$

Let  $\mathbf{z}^i=\mathbf{x}^i+\mathbf{j}\mathbf{y}^i$  ( $\mathbf{z}^i=\mathbf{x}^i-\mathbf{j}\mathbf{y}^i$ ),  $\mathbf{j}^2=1$ ,  $1\leq i\leq m$  be a para-complex local coordinate system on a neighbourhood  $\mathbf{TU}_p$  of any point  $\mathbf{T}_p$  of  $\mathbf{TM}$ . We define the vector fields as:

$$\mathbf{J}(\partial/\partial\mathbf{z}^i)_p=\frac{1}{2}\{(\partial/\partial\mathbf{x}^i)_p-\mathbf{j}(\partial/\partial\mathbf{y}^i)_p\}, \mathbf{J}(\partial/\partial\mathbf{z}^i)_p=\frac{1}{2}\{(\partial/\partial\mathbf{x}^i)_p+\mathbf{j}(\partial/\partial\mathbf{y}^i)_p\}, \quad (5)$$

and the dual co-vector fields as:

$$(\mathbf{d}\mathbf{z}^i)_p=(\mathbf{d}\mathbf{x}^i)_p+\mathbf{j}(\mathbf{d}\mathbf{y}^i)_p, (\mathbf{d}\mathbf{z}^i)_p=(\mathbf{d}\mathbf{x}^i)_p-\mathbf{j}(\mathbf{d}\mathbf{y}^i)_p, \quad (6)$$

which represent the bases of the tangent space  $T_p(\mathbf{TM})$  and cotangent space  $T_p^*(\mathbf{TM})$  of  $\mathbf{TM}$ , respectively. Then the following can be found

$$\mathbf{J}(\partial/\partial z^i) = -\mathbf{j}(\partial/\partial z^i)_p, \mathbf{J}(\partial/\partial \bar{z}^i) = \mathbf{j}(\partial/\partial \bar{z}^i). \quad (7)$$

The dual endomorphism  $\mathbf{J}^*$  of the cotangent space  $T_p^*(\mathbf{TM})$  at any point  $T_p$  of manifold  $\mathbf{TM}$  satisfies  $\mathbf{J}^{*2} = \mathbf{I}$  and is defined by

$$\mathbf{J}^*(dz^i) = -(d\bar{z}^i), \mathbf{J}^*(d\bar{z}^i) = dz^i. \quad (8)$$

### 2.2 Para-Kählerian Manifolds

**Definition 2:** An almost para-Hermitian manifold  $(\mathbf{TM}, \mathbf{g}, \mathbf{J})$  is a differentiable manifold  $\mathbf{TM}$  endowed with an almost product structure  $\mathbf{J}$  and a pseudo-Riemannian metric  $\mathbf{g}$ , compatible in the sense that

$$\mathbf{g}(\mathbf{JX}, \mathbf{Y}) + \mathbf{g}(\mathbf{X}, \mathbf{JY}) = 0, \text{ for all } \mathbf{X}, \mathbf{Y} \in \mathfrak{X}(\mathbf{TM}). \quad (9)$$

An almost para-Hermitian structure on a differentiable manifold  $\mathbf{TM}$  is  $\mathbf{G}$ -structure on  $\mathbf{TM}$  whose structural group is the representation of the para unitary group  $\mathbf{U}(\mathbf{n}, \mathbf{A})$  given in [7]. A para-Hermitian manifold is a manifold with an integrable almost para-Hermitian structure  $(\mathbf{g}, \mathbf{J})$ . 2-covariant skew tensor field  $\Phi$  defined by  $\Phi(\mathbf{X}, \mathbf{Y}) = \mathbf{g}(\mathbf{X}, \mathbf{JY})$  is called fundamental 2-form. An almost para-Hermitian manifold  $(\mathbf{TM}, \mathbf{g}, \mathbf{J})$ , such that  $\Phi$  is closed shall be called an almost para-Kählerian manifold.

A para-Hermitian manifold  $(\mathbf{TM}, \mathbf{g}, \mathbf{J})$  is said to be a para-Kählerian manifold if  $\Phi$  is closed. Also, by means of geometric structures, one may show that  $(\mathbf{T}^*\mathbf{M}, \mathbf{g}, \mathbf{J})$  is a para-Kählerian manifold.

### 2.3 Para Hamiltonian Equations

Here, we obtain para-Hamiltonian equations for classical mechanics structured on para-Kählerian manifold  $\mathbf{T}^*\mathbf{M}$ .

Let  $\mathbf{T}^*\mathbf{M}$  be any para-Kählerian manifold and  $\{z^i, \bar{z}^i\}, 1 \leq i \leq m$  its para-complex coordinates. Suppose that  $\{(\partial/\partial z^i)_p, (\partial/\partial \bar{z}^i)_p\}$  and  $\{(dz^i)_p, (d\bar{z}^i)_p\}$  are bases over the set of para-complex numbers  $\mathbf{A}$  of tangent space  $T_p(\mathbf{TM})$  and the cotangent space  $T_p^*(\mathbf{TM})$  of  $\mathbf{TM}$ . Taking care of almost para-complex structure  $\mathbf{J}^*$  given by Eq.(8), para-Liouville form  $\lambda$  is calculated as  $\lambda = \mathbf{J}^*(\omega) = 1/2 \mathbf{j}(z^i dz^i - \bar{z}^i d\bar{z}^i)$  such that para-complex 1-form  $\omega = 1/2 (z^i dz^i + \bar{z}^i d\bar{z}^i)$  on  $\mathbf{T}^*\mathbf{M}$ . If  $\Phi = -d\lambda$  is closed para-Kählerian form, then  $\Phi$  is also a para-symplectic structure on  $\mathbf{T}^*_p\mathbf{M}$ .

**Proposition 1:** Let  $T^*M$  be para-Kählerian manifold with closed para-Kählerian form  $\Phi$ . Para-Hamiltonian vector field  $Z_H$  on para-Kählerian manifold with closed para-Kählerian form  $\Phi$  is given by

$$Z_H = -j \partial H / \partial z_i \partial / \partial z_i + j \partial H / \partial z_i \partial / \partial z_i, 1 \leq i \leq m \quad (10)$$

on  $T^*M$ .

By para-Hamiltonian equations on para-Kählerian manifold  $T^*M$ , we call the following equations:

$$dz_i/dt = -j \partial H / \partial \dot{z}_i, d\dot{z}_i/dt = j \partial H / \partial z_i \quad (11)$$

### 3. Poisson Manifolds

Let  $T^*M$  be a para-Kählerian manifold with closed para-Kählerian form  $\Phi$ . If the closed para-Kählerian form  $\Phi$  on  $T^*M$  is symplectic structure, all para-Kählerian manifolds are also symplectic manifolds.

**Proposition 2:** Let  $(T^*M, \Phi)$  and  $(S, \omega)$  be symplectic manifolds of same dimension and let  $H$  be a symplectic transformation from  $(T^*M, \Phi)$  to  $(S, \omega)$ . Then  $(Th)Z_{F \circ H} = Z_F$  holds for any function  $F$  on  $S$ .

Assume that  $(T^*M, \Phi)$  is a symplectic manifold. Let  $F$  and  $G$  be  $C^\infty$  para-complex functions on  $T^*M$ . Then Poisson bracket of  $F$  and  $G$  is defined by

$$\{F, G\} = \Phi(Z_F, Z_G), \quad (12)$$

where  $Z_F, Z_G$  are Hamiltonian vector fields on  $T^*M$  defined by  $i_{Z_F} \Phi = \Phi(Z_F) = dF$  and  $i_{Z_G} \Phi = \Phi(Z_G) = dG$ , respectively.

**Definition 3:** Poisson structure is called a bilinear map defined by

$$\begin{aligned} C^\infty(T^*M) \times C^\infty(T^*M) &\rightarrow C^\infty(T^*M) \\ (F, G) &\rightarrow \{F, G\} \end{aligned} \quad (13)$$

on a para-Kählerian manifold  $T^*M$  if the following identities are verified.

- (i) (Skew symmetry)  $\{F, G\} = -\{G, F\}$ ,
- (ii) (Jacobi identity)  $\{F, \{G, H\}\} + \{G, \{H, F\}\} + \{H, \{F, G\}\} = 0$ ,
- (iii) (Leibniz rule)  $\{FG, H\} = F\{G, H\} + \{F, G\}H$ ,

where  $C^\infty(T^*M)$  is the space of  $C^\infty$  functions on  $T^*M$ .

Para-Kählerian manifold endowed with Poisson structure  $\{ , \}$  is also called Poisson manifold.

Let  $\{ , \}$  be a Poisson structure on the para-Kählerian manifold.

From (iii) of **Definition 3** we see that the map

$$\begin{aligned} \{ F, \cdot \}: C^\infty(T^*M) &\rightarrow C^\infty(T^*M), \\ (G) &\rightarrow \{ F, G \} \end{aligned} \quad (14)$$

is a derivation. Therefore there is a unique vector field  $Z_G$  on  $T^*M$  such that

$$Z_G F = \{ F, G \}, \quad (15)$$

where  $Z_G$  is said the Hamiltonian vector field of  $C^\infty$  para-complex function  $G$  on para-Kählerian manifold  $T^*M$  with closed para-Kählerian form  $\Phi$ .

#### 4. Para Hamiltonian Equations with Poisson brackets

In this section, we obtain para-complex Hamiltonian equations on  $T^*M$  with Poisson structure  $\{ , \}$ . Taking **Eq.(10)**, we deduce that the Poisson bracket of two functions  $F$  and  $G$  is

$$\{ F, G \} = -j \partial F / \partial z_i \partial G / \partial \bar{z}_i + j \partial F / \partial \bar{z}_i \partial G / \partial z_i. \quad (16)$$

From **Eq.(16)**, we obtain the Poisson brackets of the canonical coordinates:

$$\{ z_i, z_j \} = \{ \bar{z}_i, \bar{z}_j \} = 0, \quad \{ z_i, \bar{z}_j \} = \delta_i^j \quad (17)$$

Furthermore, if  $F$  is a function on  $T^*M$ , we get

$$\{ F, z_i \} = -\{ z_i, F \} = j \partial F / \partial \bar{z}_i, \quad \{ F, \bar{z}_i \} = -\{ \bar{z}_i, F \} = j \partial F / \partial z_i \quad (18)$$

Finally, using **Eqs.(11)** and **(18)**, para-complex Hamilton equations with Poisson structure are calculated as

$$dz_i/dt = -\{ z_i, H \}, \quad d\bar{z}_i/dt = -\{ \bar{z}_i, H \} \quad (19)$$

Now, we have a question. We try to solve it. A canonical transformation  $h: (T^*M, \Phi) \rightarrow (T^*M, \Phi)$  preserves para-complex Hamiltonian equations? To see this it is sufficient to show that the Poisson brackets are invariant under the action of  $H$ . Firstly, let us take

$$h^* \{ F, G \} = \{ F, G \} \circ h \quad (20)$$

Then we have

$$\{F, G\}_{oh} = (Z_G F)_{oh} = (((Th)Z_G)_{oh} F)_{oh} = (Z_G)_{oh} (F)_{oh} = \{F_{oh}, G_{oh}\} = \{h^*F, h^*G\}, \quad (21)$$

by means of Eq.(15) and Proposition 2. Hence considering Eqs.(20) and (21) it is

$$h^*\{F, G\} = \{h^*F, h^*G\} \text{ or } \{F, G\}_{oh} = \{F_{oh}, G_{oh}\}. \quad (22)$$

Finally, Poisson brackets are seen invariant under the action of  $h$ .

Especially, if  $h:(z_i, \dot{z}_i) \rightarrow (\check{z}_i, \check{z}'_i)$ , where  $(z_i, \dot{z}_i)$  and  $(\check{z}_i, \check{z}'_i)$  are canonical coordinates on  $T^*M$ , we have

$$h^*\{z_i, H\} = \{z_i, H_{oh}\} = \{\check{z}_i, K\} = d\check{z}_i/dt, \quad h^*\{\dot{z}_i, H\} = \{\dot{z}_i, H_{oh}\} = \{\check{z}'_i, K\} = d\check{z}'_i/dt, \quad (23)$$

such that  $K = h^*H = H_{oh}$ . Thus we call to be Kamiltonian  $K$  in a canonically transformed set of coordinates. It is said that  $K$  is identical to  $H$ , with the possible exception of an arbitrary additive constant if

$$K \equiv H + d\Omega/dt, \quad (24)$$

where  $\Omega$  is any function of phase space coordinates with continuous second derivatives.

### Conclusion

Taking care of the considerations the above, it is clear that Poisson bracket is the most important operation given by the symplectic and/or Kählerian structure. We conclude that the Hamiltonian formalisms in generalized classical mechanics and field theory can be intrinsically characterized on the para-Kählerian manifold endowed with Poisson structure  $\{ , \}$ . The geometric approach of para-complex Hamiltonian systems is that solutions of Hamiltonian vector field  $Z_H$  on para-Kählerian manifold  $T^*M$  are paths para-complex Hamiltonian equations obtained in Eq. (19) on  $T^*M$  with Poisson bracket  $\{ , \}$ .

With respect to Eq. (23), it was shown that canonical transformations preserve the form of Hamiltonian equations.  $\Phi$  being useful in Hamiltonian mechanics as well as thermodynamics given in Eq. (24) is known generating function for canonical transformation. Moreover, four types of time-dependent generating functions are possible to define. By means of these, it may obtain generalized Maxwell relations [8].

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